Influence of the Aspect Ratio on the Thermal Performance of Axial Extended Surface in Electric Motor by the Generalized One-Dimensional Radial Fin Using the Frobenius Method

Élcio Nogueira

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ABSTRACT: Analytical solution using the Frobenius Method for the generalized one-dimensional radial fin is applied to analyze the influence of the aspect ratio on the temperature distribution, heat transfer rate, efficiency, and effectiveness on the extended axial surface of an electric motor. Cast Iron and Aluminum materials were used for the fin. For high value for the aspect ratio and external heat transfer coefficient, the temperature, along the entire length of the fin, tends to the temperature of the medium. In contrast, the fin temperature tends to the surface temperature of the electric motor when the heat transfer coefficient decreases. The thermal performance does not increase indefinitely with the increase of the aspect ratio, and for relatively high values for aspect ratio, the heat transfer rate remains stagnant, with a tendency to decrease. The aspect ratio does not have much influence on the results in the efficiency and effectiveness, but the difference of the heat conduction coefficients between the used materials are very relevant.

KEY WORDS: Frobenius Method; Axial Extended Surface; One-dimensional Radial Fin; Electric Motor, Cast Iron; Aluminum.

I. INTRODUCTION

The serial expansion of the modified power, the Frobenius Method, is applied to a specific geometry, Figure 1, and converted to a special and particular geometry called one-dimensional radial fin in the specialized literature, Figure 2.

Extended surfaces are extensively used in engineering applications to increase heat transfer efficiency. The longitudinal fin of rectangular profile and the one-dimensional radial fin are the most commonly used Suárez[1], Soares and Nogueira[2], Santos[3], Freire et al.[4], Yuri et al.[5], Novaes et al.[6], Campo and Kundu[7], Cotta and Ramos[8], Sommers and Jacobi[9], Kraus[10], Aparecido and Cotta[11,12].

The applications in compact heat exchanger increase the interest in accessible and applicable models for fins systems Oktay[13].

Nogueira[14] demonstrated that the one-dimensional model, using Frobenius Method, is suitable for compact fins systems and that the smaller the value of the aspect ratio, the higher the range of Biot number in which the one-dimensional model works correctly.

Nogueira[15] presented a comparison between the Frobenius Method and the mathematical model already established and validated, called "Improved One-Dimensional Classic Radial Fin" in the literature Cotta and Mikhailov[16]. The equivalence of the one-dimensional models presented is evident, considering the Reynolds number used for analysis. For relatively low values of the Biot number, the efficiency and effectiveness justify the placement of the fins and use of one of the two one-dimensional models.

II. OBJECTIVES

To analyze the influence of the aspect ratio on the thermal performance of the radial fin using a generalized onedimensional model by the Frobenius method.



III. THEORETICAL ANALYSIS Generalized One-Dimensional Radial Fin by Frobenius Method

Figure 1 – Generalized Geometry for One-Dimensional Radial Fin Analysis

Basic information for the fin system	
Engine width - L_M	130,13 mm
Width of fin base - w	5,84 mm
Internal radius – r _i	69,8 mm
External radius - ro	86,8 mm
Fin height - <i>L_o</i>	17,00 mm
Conductivity - k	80W/(m.K)
Maximum base Temperature - T _b	98°C
Maximum external Temperature - T_{∞}	40°C

Table 1 – Data for typical finned Radial Electric Motor DC

Consider steady-state, one-dimensional heat conduction through a radial fin, with constant conductivity, k, and subjected an ambient temperature, T_{∞} .

$$\frac{d}{dr}\left[r\frac{dT(r)}{dr}\right] = \frac{2h_2[\alpha r + L_M]}{k\alpha L_M}[T(r) - T_\infty]$$
(01)

By definition:

$$\theta(r) = \frac{T(r) - T_{\infty}}{T(r_i) - T_{\infty}} \tag{02}$$

Then, we have:

$$r^{2}\frac{d^{2}\theta(r)}{dr^{2}} + r\frac{d\theta(r)}{dr} = \frac{2h_{2}[\alpha r^{2} + L_{M}r]}{k\alpha L_{M}}\theta(R)$$
(03)

$$R = \frac{r - r_i}{r_o - r_i} \quad \rightarrow \quad r = L_o R + r_i \text{ with } L_o = r_o - r_i \tag{04}$$

$$[R + \frac{r_i}{L_o}]^2 \frac{d^2 \theta(R)}{dR^2} + \left(R + \frac{r_i}{L_o}\right) \frac{d\theta(R)}{dR} = \frac{2h_2[\alpha(L_oR + r_i)^2 + L_M(L_oR + r_i)]}{k\alpha L_M} \theta(R)$$
(05)
$$\mathbb{P}(R) \frac{d^2 \theta(R)}{dR^2} + \mathbb{Q}(R) \frac{d\theta(R)}{dR} - \mathbb{W}(R) \theta(R) = 0$$
(06)

At where

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$$\mathbb{P}(R) = P_1 R^2 + P_2 R + P_3; \ \mathbb{Q}(R) = Q_1 R + Q_2; \ \mathbb{W}(R) = W_1 R^2 + W_2 R + W_3$$
(07)

$$P_{1} = 1.0; P_{2} = 2K_{1}; P_{3} = K_{1}^{2}; Q_{1} = 1.0; Q_{2} = K_{1}$$

$$2R_{2}K^{2} \qquad [4K^{2}K_{2} - 2K] \qquad 2K^{2}K_{2} - 2KK_{2}$$
(08)

$$W_1 = \frac{2B_{i2}K^2}{K^2}; W_2$$

and

$$_{1} = \frac{2B_{i2}K^{2}}{K_{2}}; W_{2} = B_{i2}\left[\frac{4K^{2}K_{1}}{K_{2}} + \frac{2K}{\alpha}\right]; W_{3} = B_{i2}\left[\frac{2K^{2}K_{1}}{K_{2}} + \frac{2KK_{1}}{\alpha}\right]$$
(09)

The dimensionless groups were defined for the problem in analysis:

$$K = \frac{L_o}{w/2}; K_1 = \frac{r_i}{L_o}; K_2 = \frac{L_M}{w/2}; B_{i1} = \frac{h_2(\frac{w}{2})}{k}; B_{i2} = \frac{h_2(\frac{w}{2})}{k}$$
(10)

For the situation under analysis, an electric motor with axial fins, radial profile, and prescribed temperature in the base, Figure 2, the following simplifications are necessary:



Figure 2 – Geometric Representation of Classical One-dimensional Radial Fin

Then, $P_2 = 0; P_3 = 0; Q_2 = 0; W_2 = 0 and W_3 = 0$ (12)And

$$R^{2} \frac{d^{2} \theta(R)}{dR^{2}} + R \frac{d\theta(R)}{dR} - W_{1} R^{2} \theta(R) = 0$$
For convenience, was defined
(13)

$$K_2 = 2; \ \beta^2 = W_1 \ and \ R' = \beta R$$
 (14)

In this case,

$$R'^{2} \frac{d^{2} \theta(R')}{dR'^{2}} + R' \frac{d\theta(R')}{dR'} - {R'}^{2} \theta(R') = 0$$
(15)

The Equation (15) has a singular regular point in R'= 0 and By Georg Frobenius (1849-1917), Euler-Mascheroni [22, pag.243], Kreyzig [17, pag.190], Arpaci [18, pag.231], Hildebrand [19, pag.143], Schneider [20, pag.46-59], Carslaw and Jaeger[21, pag.374-376]:

$$\theta(R') = \sum_{n=0}^{\infty} a_n {R'}^{n+s}$$
(16)

$$\theta'(R') = \frac{d\theta(R')}{dR'} = \sum_{n=1}^{\infty} a_{n-1}(n+s-1){R'}^{n+s}$$
(17)

$$\theta''(R') = \frac{d^2\theta(R')}{d{R'}^2} = \sum_{n=2}^{\infty} a_{n-2}(n+s-2)(n+s-3){R'}^{n+s}$$
(18)

Then

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(11)

$$\mathbb{P}(R)\sum_{n=2}^{\infty}a_{n-2}(n+s-2)(n+s-3)R^{n+s} + \mathbb{Q}(R)\sum_{n=1}^{\infty}a_{n-1}(n+s-1)R^{n+s} - \mathbb{W}(R)\sum_{n=0}^{\infty}a_nR^{n+s} = 0$$
(19)

n=0 By algebraic manipulation, the following indicial equation was obtained:

$$a_0[(s^2 - s) + s]R'^s = 0 \quad with \quad a_0 \neq 0 \text{ and } s = 0$$
(20)
The roots of the indicial equation are equal to zero, and the recurrence rule is given by

$$a_n = \frac{a_{n-2}}{n^2} \tag{21}$$

or

$$a_2 = \frac{a_0}{2^2}; \ a_4 = \frac{a_0}{2^2 4^2}; \ a_6 = \frac{a_0}{2^2 4^2 6^2} \dots$$
 (22)

For the situation in analysis, two equal roots, there are two linearly independent solutions, which constitute a fundamental system of solution Kreyzig [17]. The first is:

$$\theta_1(R) = 1 + \sum_{m=1,2,3..}^{\infty} a_{2m} (\beta R)^{2m}; \quad a_{2m} = \frac{1}{2^{2m} (m!)^2}$$
(23)

The second linearly independent solution contains a logarithmic term and has a form:

$$\theta_2(R) = [\ln(\beta R)]\theta_1(R) + \sum_{m=1,2,3...} A_m(\beta R)^m$$
(24)

By Carslaw and Jaeger[21], and Boyce and Diprima[22] the more convenient expression is

$$\theta_2(R) = -\left[\ln\left(\frac{\beta R}{2}\right) + \gamma\right]\theta_1(R) + \sum_{m=1,2,3\dots}^{\infty} a_{2m} H_m(\beta R)^{2m}$$
(25)

At where

$$H_m = \frac{1}{m} + \frac{1}{m-1} + \dots + \frac{1}{2} + 1 \quad and \ \gamma \cong 0.5772 \tag{26}$$

n as the Euler-Mascheroni [22, pag 247] constant.

 γ is known as the Euler-Mascheroni[22, pag.247] constant. Then

$$\theta(R) = a_0 \theta_1(R) + a_1 \theta_2(R)$$
(27)
$$\theta(R) = a_0 [1 + \sum_{m=1,2,3...}^{\infty} a_{2m} (\beta R)^{2m}] + a_1 \left[-\left[\ln\left(\frac{\beta R}{2}\right) + \gamma \right] \theta_1(R) + \sum_{m=1,2,3...}^{\infty} a_{2m} H_m (\beta R)^{2m} \right]$$
(28)

The first boundary condition is defined by Cotta and Mikhailov[16]:

$$\theta(0) = 1 \rightarrow a_0 = 1 + a_1 \left[\ln \left(\frac{\beta R_b}{2} \right) + \gamma \right]$$
(29)

Finally,

$$\theta(R) = \theta_1(R) + a_1 \left\{ \left[\ln\left(\frac{\beta R_b}{2}\right) + \gamma \right] \theta_1(R) + \theta_2(R) \right\}$$
(30)

$$\theta'(R) = \theta_1'(R) + a_1 \left\{ \left[\ln\left(\frac{\beta R_b}{2}\right) + \gamma \right] \theta_1'(R) + \theta_2'(R) \right\}$$
(31)

Where (Figure 2)

$$R_b = R \to 0$$
For the second boundary conditions: (32)

$$\theta'(1) = -B_{i2}K\theta(1) \tag{33}$$

In this case,

$$\frac{-[\theta_1(1) + B_{i2}K\theta_1'(1)]}{(34)}$$

$$a_{1} = \frac{\left[1 + B_{12}K\theta_{1}(1)\right]}{\left[\ln\left(\frac{\beta R_{b}}{2}\right) + \gamma\right] \left[\theta_{1}(1) + B_{i2}K\theta_{1}'(1)\right] + \left[\theta_{2}(1) + B_{i2}K\theta_{2}'(1)\right]}$$
(5)

$$\dot{q} = \frac{-kA_b(T_b - T_{\infty})\theta'(0)}{L_0}$$
(35)

The dimensionless exchange heat transfer is written in the form, by definition

$$Q_b = \frac{q}{h_2 A_b (T_b - T_\infty)} \longrightarrow Q_b = \frac{-1}{B_{i2} K} (\frac{d\theta}{dR})_{R=0}$$
(36)

 A_b and T_b are the base area and the base temperature respectively Efficiency is given by

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$$\eta = \frac{-1}{B_{i2}K(1+K)} \left(\frac{d\theta}{dR}\right)_{R=0}$$
(37)

And, the efficacy

$$\varepsilon = \frac{2Q_b}{B_{i2}K}$$

IV. RESULTS AND DISCUSSION

Results for temperature profile, heat transfer rate, efficiency, and effectiveness were obtained for different values of aspect ratio ($K = \frac{L_0}{w/2}$) and two different materials used in electric motors (Cast Iron and Aluminum).

The parameter R_b (equations 29 – 31) has a great influence on the generation of the results obtained by the model and, in general, it is strongly dependent on the heat transfer coefficient (h₂) and the aspect ratio. Figure 3, below, presents for the high amplitude of the heat transfer coefficient and a wide range of values for the aspect ratio, the relationship of R_b to the heat transfer coefficient external to the motor surface.

Figure 4 shows temperature distribution as a function of the radial position, for a wide range of values of the external heat transfer coefficient and high value for the aspect ratio. For a high value of the heat transfer coefficient by external convection ($h_2 = 1000$), the temperature, along the entire length of the fin, tends to the temperature of the medium. In contrast, the fin temperature tends to the surface temperature of the electric motor when the heat transfer coefficient decreases ($h_2 = 50$).



FIGURE 3- Rb versus convection heat transfer coefficient h2

(38)



FIGURE 4–Dimensional temperature versus radial fin position



FIGURE 5–Dimensional heat transfer rate versus heat transfer coefficient h₂

The dimensionless heat transfer rate, as a function of the heat transfer coefficient by external convection, is represented by Figure 5, for different values of the aspect ratio and Cast Iron as the material of the fin. The results demonstrate that the thermal performance does not increase indefinitely with the increase of the aspect ratio. For relatively high values for aspect ratio, the heat transfer rate remains stagnant, with a tendency to decrease. This fact can be justified by referring to Figure 4 when the temperature tends to the value of the surface temperature of the electric motor for high values of the aspect ratio and high values for the external heat transfer coefficient.

Figure 6 presents results for efficiency as a function of the external heat transfer coefficient for two types of materials for the fin of the electric motor, Cast Iron, and Aluminum. For Cast Iron, efficiency is justified, above 60%, for values of the external heat transfer coefficient below 150 W / (m^2 . K). On the other hand, for Aluminum, values below 400 W / (m^2 . K) enable efficiency above 60% for the fin of the electric motor. The fact to be emphasized, concerning the efficiency, is that the aspect ratio is relevant for the high external heat transfer coefficient. For low values of the external heat transfer coefficient, increasing the aspect

International organization of Scientific Research

ratio does not have a significant influence on the efficiency. The difference, observable, concerning Cast Iron and Aluminum, can be justified according to the values of the conduction heat transfer coefficient of the materials: k = 80 W / (m. K) for Cast Iron and k = 237 W / (m. K) for Aluminum.

Values for efficacy, considering Cast Iron and Aluminum, are shown in Figure 7. The aspect ratio does not have much influence on the results, and the difference between the heat conduction coefficients is very relevant. To efficacy, as values below 2 do not justify the placement of fins, values above 200 W / (m^2 K) are not recommended for cast iron, and values above 400 W / (m^2 K) is not recommended for Aluminum.



FIGURE 6 – Effectiveness versus heat transfer coefficient h2



FIGURE 7 -Efficacy versus heat transfer coefficient h₂

V. CONCLUSIONS

The main conclusions reached in this work are described below:

- The thermal performance does not increase indefinitely with the increase of the aspect ratio.

- For relatively high values for aspect ratio, the heat transfer rate remains stagnant, with a tendency to decrease.

- Regarding efficiency, the aspect ratio is only relevant for high values of the external heat transfer coefficient.

- For Cast Iron, efficiency is justified for values of the external heat transfer coefficient below 200 W / (m^2, K) .

- For Aluminum, values below 400 W / (m 2 . K) enable efficiency above 60% for the fin of the electric motor.

- The aspect ratio does not have much influence on the results of the efficacy, but the difference between the heat conduction coefficients, Cast Iron and Aluminum, is very relevant.

The synthesis of the results obtained lead us to conclude that the aspect ratio is relevant for values less than and equal to 9, and for values of the external heat transfer coefficient less than 200 W / (m^2 K) for Cast Iron, and less than 400 W / (m^2 K) for Aluminum.

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